

A survey on Piecewise Linear Iterated Function Systems on the line

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Introduction

The main object of our investigation is a special family of iterated function systems, defined on the real line.

Definition 1.1

If $\mathcal{F} = \{f_k\}_{k=1}^m$ is a finite list of strict contractions on \mathbb{R} , then we call \mathcal{F} an iterated function system (IFS).

Each Iterated function system \mathcal{F} uniquely defines an invariant set that we refer to as the attractor of \mathcal{F} .

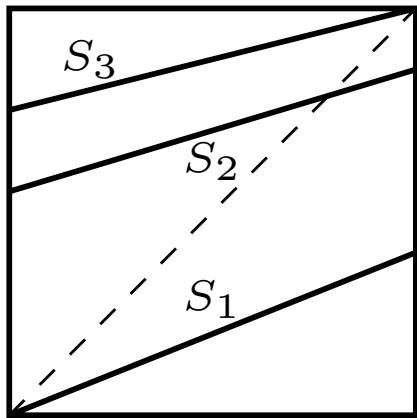
Definition 1.2

The **attractor** $\Lambda^{\mathcal{F}}$ of the IFS \mathcal{F} is the unique non-empty compact set that satisfies

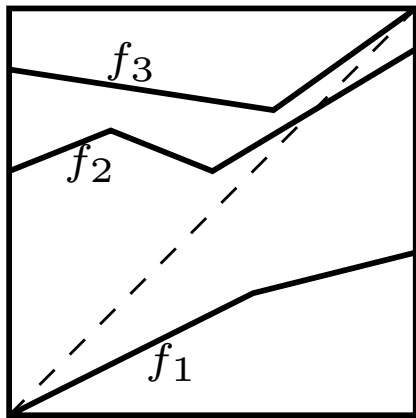
$$(1) \quad \Lambda^{\mathcal{F}} = \bigcup_{k=1}^m f_k(\Lambda^{\mathcal{F}}).$$

When all the f_i functions in \mathcal{F} are linear, then \mathcal{F} is called **self-similar iterated function system**. A lot is known about these systems, but overlapping constructions can be hard to handle.

The functions in our systems are only **piecewise linear**. Each function f_i might change slope at some point, thus they are not necessarily injective. Hence, overlapping structures are even more typical in our case.



Self-similar IFS
 $\mathcal{S} = \{S_1, S_2, S_3\}$



CPLIFS
 $\mathcal{F} = \{f_1, f_2, f_3\}$

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One of the central questions in the field of fractal geometry is the dimension of the attractor of iterated function systems.

Definition 2.1

The Hausdorff dimension of a set $A \subset \mathbb{R}$ is defined as

$$\dim_{\text{H}} A = \inf \left\{ s \geq 0 \mid \forall \varepsilon > 0, \exists \{U_i\}_{i=1}^{\infty} : A \subset \bigcup_{i=1}^{\infty} U_i \text{ \& } \sum_{i=1}^{\infty} |U_i|^s \leq \varepsilon \right\}.$$

The definition implies that the dimension of the attractor is connected to the yellow covering sum.

Without loss of generality, assume that the attractor $\Lambda^{\mathcal{F}}$ is contained in $[0, 1]$, then by iteration

$$(2) \quad \Lambda^{\mathcal{F}} = \bigcap_{n=1}^{\infty} \bigcup_{(i_1, \dots, i_n) \in [m]^n} f_{i_1 \dots i_n}([0, 1]).$$

The sets $f_{i_1 \dots i_n}([0, 1])$ are called cylinder intervals, and they are the most natural cover of the attractor. We use these sets to define the natural pressure function of \mathcal{F}

$$(3) \quad \Phi^{\mathcal{F}}(s) := \limsup_{n \rightarrow \infty} \frac{1}{n} \log \sum_{i_1 \dots i_n} |f_{i_1 \dots i_n}([0, 1])|^s$$

According to Barreira [1] the function $\Phi^{\mathcal{F}}(s)$ is strictly decreasing, continuous, positive at 0 and tends to $-\infty$ as $s \rightarrow \infty$. Thus we can define the **natural dimension** of \mathcal{F} as

$$(4) \quad s_{\mathcal{F}} := (\Phi^{\mathcal{F}})^{-1}(0).$$

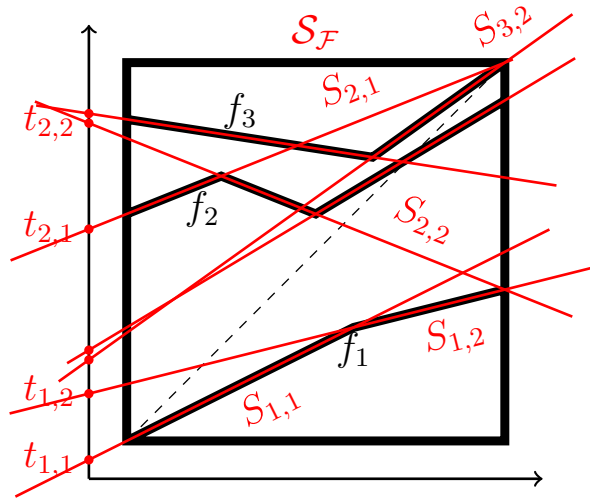
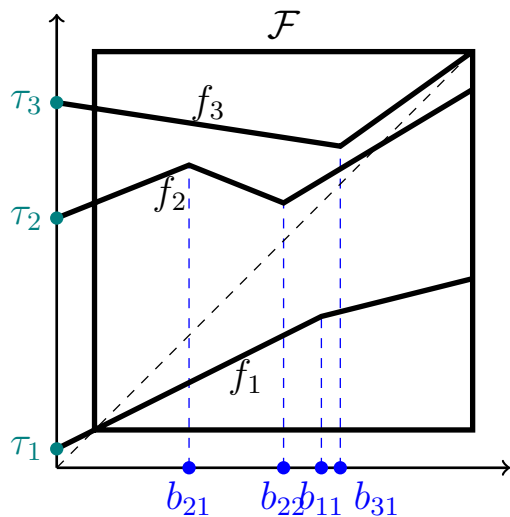
We showed that for the attractor of a typical CPLIFS the natural dimension equals to the Hausdorff dimension.

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Definition 3.1

We say that a small CPLIFS \mathcal{F} is **regular** if its attractor $\Lambda_{\mathcal{F}}$ does not contain any of the breaking points.

For a regular CPLIFS \mathcal{F} there is a smallest **N** such that

(5) there are no breaking points in $\bigcup_{i_1 \dots i_N} f_{i_1 \dots i_N}([0, 1])$.

Theorem 3.2 (P., Simon [2])

For a $\dim_{\mathbf{P}}$ -typical small CPLIFS \mathcal{F} we have

$$(6) \quad \dim_{\mathbf{H}} \Lambda^{\mathcal{F}} = \dim_{\mathbf{B}} \Lambda^{\mathcal{F}} = s_{\mathcal{F}}.$$

Theorem 3.3 (P., Simon [2])

*Let \mathcal{F} be a **regular** CPLIFS for which the **generated self-similar IFS** satisfies the **Exponential Separation Condition (ESC)**. Then*

$$(7) \quad \dim_{\mathbf{H}} \Lambda^{\mathcal{F}} = \dim_{\mathbf{B}} \Lambda^{\mathcal{F}} = s_{\mathcal{F}}.$$

Thank you for your attention!

References

[1] Luis M Barreira.

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Ergodic Theory and Dynamical Systems, 16(5):871–928, 1996.

[2] R Dániel Prokaj and Károly Simon.

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